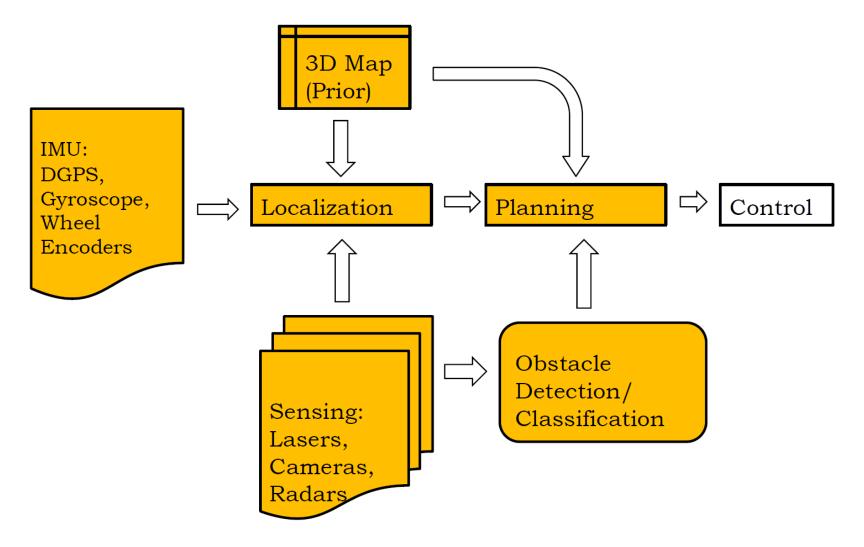
Attitude Representation and Transformation Matrices

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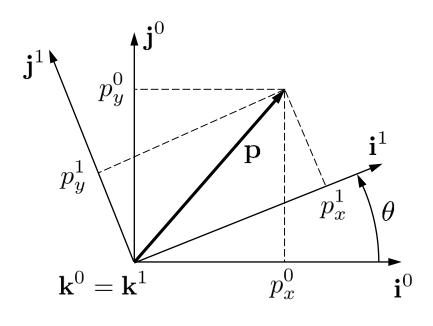
Autonomous Navigation System



Coordinated Frames

- Describe relative position and orientation of objects
 - Aircraft relative to direction of wind
 - Camera relative to aircraft
 - Aircraft relative to inertial frame
- Some things most easily calculated or described in certain reference frames
 - Newton's law
 - Aircraft attitude
 - Aerodynamic forces/torques
 - Accelerometers, rate gyros
 - GPS
 - Mission requirements

Rotation of Reference Frame



$$\mathbf{p} = p_x^0 \mathbf{i}^0 + p_y^0 \mathbf{j}^0 + p_z^0 \mathbf{k}^0$$
$$\mathbf{p} = p_x^1 \mathbf{i}^1 + p_y^1 \mathbf{j}^1 + p_z^1 \mathbf{k}^1$$
$$p_x^1 \mathbf{i}^1 + p_y^1 \mathbf{j}^1 + p_z^1 \mathbf{k}^1 = p_x^0 \mathbf{i}^0 + p_y^0 \mathbf{j}^0 + p_z^0 \mathbf{k}^0$$
$$\mathbf{p}^1 \stackrel{\triangle}{=} \begin{pmatrix} p_x^1 \\ p_y^1 \\ p_z^1 \end{pmatrix} = \begin{pmatrix} \mathbf{i}^1 \cdot \mathbf{i}^0 & \mathbf{i}^1 \cdot \mathbf{j}^0 & \mathbf{i}^1 \cdot \mathbf{k}^0 \\ \mathbf{j}^1 \cdot \mathbf{i}^0 & \mathbf{j}^1 \cdot \mathbf{j}^0 & \mathbf{j}^1 \cdot \mathbf{k}^0 \\ \mathbf{k}^1 \cdot \mathbf{i}^0 & \mathbf{k}^1 \cdot \mathbf{j}^0 & \mathbf{k}^1 \cdot \mathbf{k}^0 \end{pmatrix} \begin{pmatrix} p_x^0 \\ p_y^0 \\ p_z^0 \end{pmatrix}$$

$$\mathbf{p}^{1} = \mathcal{R}_{0}^{1} \mathbf{p}^{0} \quad \text{where} \quad \mathcal{R}_{0}^{1} \stackrel{\triangle}{=} \begin{pmatrix} \cos \theta & \sin \theta & 0 \\ -\sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{pmatrix}$$
(rotation about **k** axis)

Rotation of Reference Frame

Right-handed rotation about \mathbf{j} axis:

$$\mathcal{R}_0^1 \stackrel{\triangle}{=} \begin{pmatrix} \cos\theta & 0 & -\sin\theta \\ 0 & 1 & 0 \\ \sin\theta & 0 & \cos\theta \end{pmatrix}$$

Right-handed rotation about **i** axis:

$$\mathcal{R}_0^1 \stackrel{\triangle}{=} \begin{pmatrix} 1 & 0 & 0\\ 0 & \cos\theta & \sin\theta\\ 0 & -\sin\theta & \cos\theta \end{pmatrix}$$

Orthonormal matrix properties:

P.1. $(\mathcal{R}_a^b)^{-1} = (\mathcal{R}_a^b)^\top = \mathcal{R}_b^a$ **P.2.** $\mathcal{R}_b^c \mathcal{R}_a^b = \mathcal{R}_a^c$ **P.3.** det $(\mathcal{R}_a^b) = 1$

Particle/Rigid Body Rotation

 One can say then that a Rigid Body is essentially a Reference Frame (RF). The translation of the origin of the RF describes the translational position. The specific orientation of the axes wrt to a chosen Inertial Reference provides the angular position.

$$\{\hat{n}\} = \begin{cases} \hat{n}_{1} \\ \hat{n}_{2} \\ \hat{n}_{2} \end{cases} \quad \{\hat{b}\} = \begin{cases} \hat{b}_{1} \\ \hat{b}_{2} \\ \hat{b}_{2} \end{cases}$$

$$\{\hat{b}\} = \begin{bmatrix} \cos \alpha_{11} & \cos \alpha_{12} & \cos \alpha_{13} \\ \cos \alpha_{21} & \cos \alpha_{22} & \cos \alpha_{23} \\ \cos \alpha_{31} & \cos \alpha_{32} & \cos \alpha_{33} \end{bmatrix} \{\hat{n}\} = [C]\{\hat{n}\}$$

$$\hat{b}_{3}$$

$$\begin{bmatrix} \mathbf{C} \end{bmatrix} - \text{Direction Cosine Matrix}$$

$$C_{ij} = \cos(\angle \hat{b}_{i}, \hat{n}_{j}) = \hat{b}_{i} \cdot \hat{n}_{j}$$

$$\hat{h}_{1}$$

$$\{\hat{n}\} - \text{Inertial Reference}$$

$$\{\hat{n}\} - \text{Body Reference}$$

$$\hat{b}_{2}$$

$$\hat{b}_{3}$$

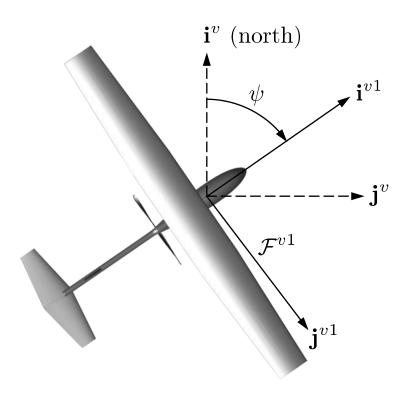
$$\hat$$

Euler Angles

- Need way to describe attitude of aircraft
- Common approach: Euler angles
 - ψ : heading (yaw)
 - θ : elevation (pitch)
 - ϕ : bank (roll)

- Pro: Intuitive
- Con: Mathematical singularity
 - Quaternions are alternative for overcoming singularity

Vehicle-1 Frame

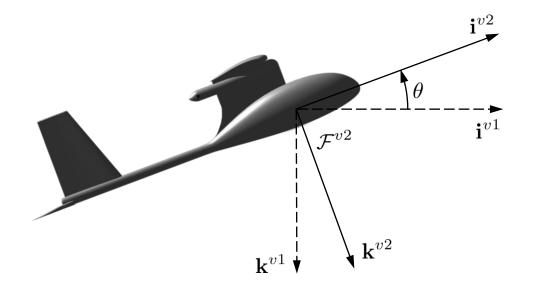


$$\mathcal{R}_v^{v1}(\psi) = \begin{pmatrix} \cos\psi & \sin\psi & 0\\ -\sin\psi & \cos\psi & 0\\ 0 & 0 & 1 \end{pmatrix}$$

$$\mathbf{p}^{v1} = \mathcal{R}_v^{v1}(\psi)\mathbf{p}^v$$

 $\psi :$ heading

Vehicle-2 Frame

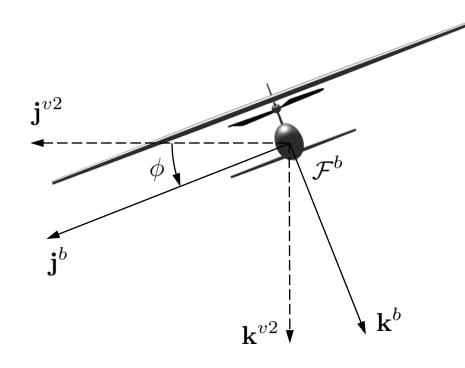


$$\mathcal{R}_{v1}^{v2}(\theta) = \begin{pmatrix} \cos\theta & 0 & -\sin\theta \\ 0 & 1 & 0 \\ \sin\theta & 0 & \cos\theta \end{pmatrix}$$

$$\mathbf{p}^{v2} = \mathcal{R}_{v1}^{v2}(\theta)\mathbf{p}^{v1}$$

 θ : elevation (pitch)

Body Frame



$$\mathcal{R}_{v2}^{b}(\phi) = \begin{pmatrix} 1 & 0 & 0\\ 0 & \cos\phi & \sin\phi\\ 0 & -\sin\phi & \cos\phi \end{pmatrix}$$

$$\mathbf{p}^b = \mathcal{R}^b_{v2}(\phi)\mathbf{p}^{v2}$$

 ϕ : bank (roll)

Inertial Frame to Body Frame Transformation

$$\begin{aligned} \mathcal{R}_{v}^{b}(\phi,\theta,\psi) &= \mathcal{R}_{v2}^{b}(\phi)\mathcal{R}_{v1}^{v2}(\theta)\mathcal{R}_{v}^{v1}(\psi) \\ &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos\phi & \sin\phi \\ 0 & -\sin\phi & \cos\phi \end{pmatrix} \begin{pmatrix} \cos\theta & 0 & -\sin\theta \\ 0 & 1 & 0 \\ \sin\theta & 0 & \cos\theta \end{pmatrix} \begin{pmatrix} \cos\psi & \sin\psi & 0 \\ -\sin\psi & \cos\psi & 0 \\ 0 & 0 & 1 \end{pmatrix} \\ &= \begin{pmatrix} c_{\theta}c_{\psi} & c_{\theta}s_{\psi} & -s_{\theta} \\ s_{\phi}s_{\theta}c_{\psi} - c_{\phi}s_{\psi} & s_{\phi}s_{\theta}s_{\psi} + c_{\phi}c_{\psi} & s_{\phi}c_{\theta} \\ c_{\phi}s_{\theta}c_{\psi} + s_{\phi}s_{\psi} & c_{\phi}s_{\theta}s_{\psi} - s_{\phi}c_{\psi} & c_{\phi}c_{\theta} \end{pmatrix} \end{aligned}$$

$$\mathbf{p}^b = \mathcal{R}^b_v(\theta)\mathbf{p}^v$$

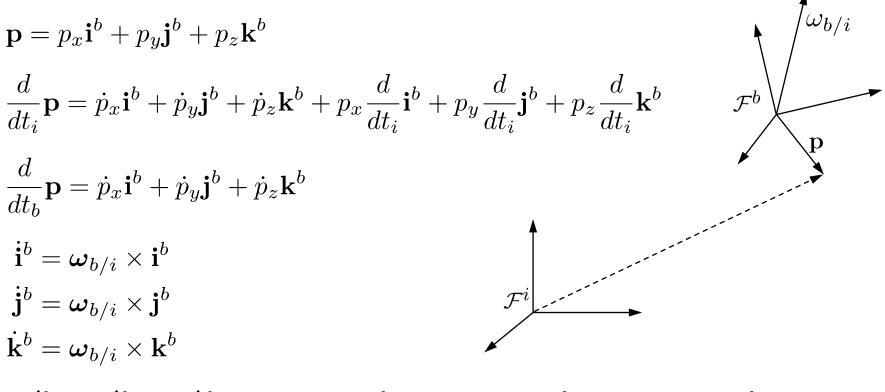
Rotational Kinematics

$$\begin{pmatrix} p \\ q \\ r \end{pmatrix} = \begin{pmatrix} \dot{\phi} \\ 0 \\ 0 \end{pmatrix} + \mathcal{R}_{v2}^{b}(\phi) \begin{pmatrix} 0 \\ \dot{\theta} \\ 0 \end{pmatrix} + \mathcal{R}_{v2}^{b}(\phi) \mathcal{R}_{v1}^{v2}(\theta) \begin{pmatrix} 0 \\ 0 \\ \dot{\psi} \end{pmatrix}$$

$$= \begin{pmatrix} \dot{\phi} \\ 0 \\ 0 \end{pmatrix} + \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \phi & \sin \phi \\ 0 & -\sin \phi & \cos \phi \end{pmatrix} \begin{pmatrix} 0 \\ \dot{\theta} \\ \dot{\theta} \end{pmatrix} + \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \phi & \sin \phi \\ 0 & -\sin \phi & \cos \phi \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ \dot{\theta} \\ \dot{\psi} \end{pmatrix}$$

$$= \begin{pmatrix} 1 & 0 & -\sin \theta \\ 0 & \cos \phi & \sin \phi \cos \theta \\ 0 & -\sin \phi & \cos \phi \cos \theta \end{pmatrix} \begin{pmatrix} \dot{\phi} \\ \dot{\theta} \\ \dot{\psi} \end{pmatrix}$$
Inverting gives
$$\begin{pmatrix} \dot{\phi} \\ \dot{\theta} \\ \dot{\psi} \end{pmatrix} = \begin{pmatrix} 1 & \sin \phi \tan \theta & \cos \phi \tan \theta \\ 0 & \cos \phi & -\sin \phi \\ 0 & \sin \phi \sec \theta & \cos \phi \sec \theta \end{pmatrix} \begin{pmatrix} p \\ q \\ r \end{pmatrix}$$

Differentiation of a Vector



 $p_x \mathbf{i}^b + p_y \mathbf{j}^b + p_z \mathbf{k}^b = p_x (\boldsymbol{\omega}_{b/i} \times \mathbf{i}^b) + p_y (\boldsymbol{\omega}_{b/i} \times \mathbf{j}^b) + p_z (\boldsymbol{\omega}_{b/i} \times \mathbf{k}^b)$ $= \boldsymbol{\omega}_{b/i} \times \mathbf{p}$

$$\frac{d}{dt_i}\mathbf{p} = \frac{d}{dt_b}\mathbf{p} + \boldsymbol{\omega}_{b/i} \times \mathbf{p}$$

Let $\{b\}$ have an angular velocity ω and be expressed as:

$$\omega = \omega_1 \hat{b}_1 + \omega_2 \hat{b}_2 + \omega_3 \hat{b}_3$$

Then
$$\frac{\mathcal{N}_{d}}{dt}\{\hat{b}_{i}\} = \frac{\mathcal{B}_{d}}{dt}\{\hat{b}_{i}\} + \omega \times \{\hat{b}_{i}\}$$

Thus $\begin{bmatrix} \frac{\mathcal{N}_{d}}{dt}\{\hat{b}\} = -[\tilde{\omega}]\{\hat{b}\} \\ \begin{bmatrix} \tilde{x}\\ \end{bmatrix} = \begin{bmatrix} 0 & -x_{3} & x_{2}\\ x_{3} & 0 & -x_{1}\\ -x_{2} & x_{1} & 0 \end{bmatrix}$
But LHS $\frac{\mathcal{N}_{d}}{dt}([C]\{\hat{n}\}) = \frac{d}{dt}([C])\{\hat{n}\} + [C]\frac{\mathcal{N}_{d}}{dt}(\{\hat{n}\}) = [\dot{C}]\{\hat{n}\}$
Finally $([\dot{C}] + [\tilde{\omega}][C])\{\hat{n}\} = 0$
Nine parameter attitude $[\dot{C}] = -[\tilde{\omega}][C]$

representation

Poisson Kinematic Equation

For the Euler 3-2-1 Sequence

$$\begin{bmatrix} C \end{bmatrix} = \begin{bmatrix} c\theta_2 c\theta_1 & c\theta_2 s\theta_1 & -s\theta_2 \\ s\theta_3 s\theta_2 c\theta_1 - c\theta_3 s\theta_1 & s\theta_3 s\theta_2 s\theta_1 + c\theta_3 c\theta_1 & s\theta_3 c\theta_2 \\ c\theta_3 s\theta_2 c\theta_1 + s\theta_3 s\theta_1 & c\theta_3 s\theta_2 s\theta_1 - s\theta_3 c\theta_1 & c\theta_3 c\theta_2 \end{bmatrix}$$

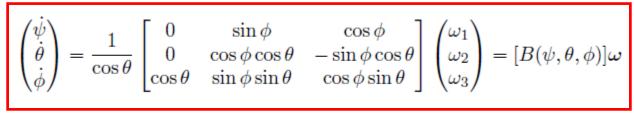
$$\psi = \theta_1 = \tan^{-1} \left(\frac{C_{12}}{C_{11}}\right)$$

$$\theta = \theta_2 = -\sin^{-1} (C_{13})$$

$$\phi = \theta_3 = \tan^{-1} \left(\frac{C_{23}}{C_{33}}\right)$$

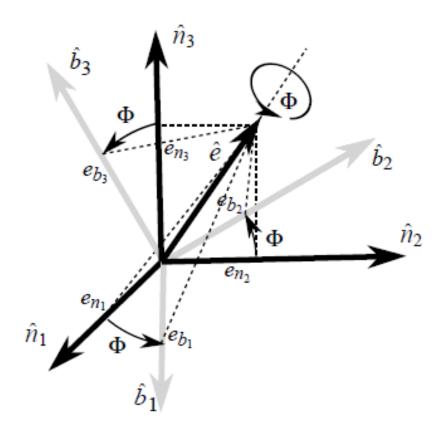
$$\begin{pmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \end{pmatrix} = \begin{bmatrix} -\sin\theta & 0 & 1 \\ \sin\phi\cos\theta & \cos\phi & 0 \\ \cos\phi\cos\theta & -\sin\phi & 0 \end{bmatrix} \begin{pmatrix} \dot{\psi} \\ \dot{\theta} \\ \dot{\phi} \end{pmatrix}$$

Attitude Kinematics Differential Equation

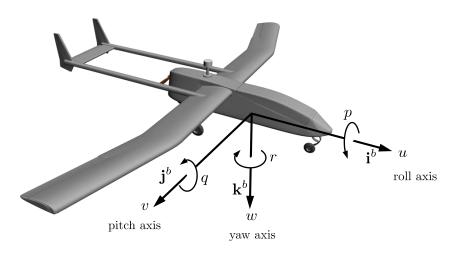


Euler's Principal Rotation Theorem

Informal Statement: There exists a principal axis about which a single axis rotation through Φ will orient the Inertial axes with the Body axes.



Rotational Dynamics



Newton's 2nd Law:

$$\frac{d\mathbf{h}}{dt_i} = \mathbf{m}$$

- h is the angular momentum vector
- m is the sum of all external moments
- Time derivative taken wrt inertial frame

$$\frac{d\mathbf{h}}{dt_i} = \frac{d\mathbf{h}}{dt_b} + \boldsymbol{\omega}_{b/i} \times \mathbf{h} = \mathbf{m}$$

Expressed in the body frame,

$$\frac{d\mathbf{h}^b}{dt_b} + \boldsymbol{\omega}^b_{b/i} \times \mathbf{h}^b = \mathbf{m}^b$$

Rotational Dynamics

$$\frac{d\mathbf{h}^b}{dt_b} + \boldsymbol{\omega}^b_{b/i} \times \mathbf{h}^b = \mathbf{m}^b$$

Because J is unchanging in the body frame, $\frac{d\mathbf{J}}{dt_b} = 0$ and $\mathbf{J} \frac{d\boldsymbol{\omega}_{b/i}^b}{dt_b} + \boldsymbol{\omega}_{b/i}^b \times (\mathbf{J} \boldsymbol{\omega}_{b/i}^b) = \mathbf{m}^b$

Rearranging we get

$$\dot{\boldsymbol{\omega}}_{b/i}^b = \mathbf{J}^{-1} \left[-\boldsymbol{\omega}_{b/i}^b \times \left(\mathbf{J} \boldsymbol{\omega}_{b/i}^b \right) + \mathbf{m}^b
ight]$$

where
$$\dot{\omega}_{b/i}^b = \frac{d\omega_{b/i}^b}{dt_b} = \begin{pmatrix} \dot{p} \\ \dot{q} \\ \dot{r} \end{pmatrix}$$

Inertia matrix

$$\mathbf{J} = \begin{pmatrix} \int (y^2 + z^2) \, d\mathbf{m} & -\int xy \, d\mathbf{m} & -\int xz \, d\mathbf{m} \\ -\int xy \, d\mathbf{m} & \int (x^2 + z^2) \, d\mathbf{m} & -\int yz \, d\mathbf{m} \\ -\int xz \, d\mathbf{m} & -\int yz \, d\mathbf{m} & \int (x^2 + y^2) \, d\mathbf{m} \end{pmatrix}$$

Rotational Dynamics

If the aircraft is symmetric about the $\mathbf{i}^b \cdot \mathbf{k}^b$ plane, then $J_{xy} = J_{yz} = 0$ and

$$\mathbf{J} = \begin{pmatrix} J_x & 0 & -J_{xz} \\ 0 & J_y & 0 \\ -J_{xz} & 0 & J_z \end{pmatrix}$$

This symmetry assumption helps simplify the analysis. The inverse of ${\bf J}$ becomes

$$\mathbf{J}^{-1} = \frac{\operatorname{adj}(\mathbf{J})}{\operatorname{det}(\mathbf{J})} = \frac{\begin{pmatrix} J_y J_z & 0 & J_y J_{xz} \\ 0 & J_x J_z - J_{xz}^2 & 0 \\ J_{xz} J_y & 0 & J_x J_y \end{pmatrix}}{J_x J_y J_z - J_{xz}^2 J_y}$$
$$= \begin{pmatrix} \frac{J_z}{\Gamma} & 0 & \frac{J_{xz}}{\Gamma} \\ 0 & \frac{1}{J_y} & 0 \\ \frac{J_{xz}}{\Gamma} & 0 & \frac{J_x}{\Gamma} \end{pmatrix} \qquad \Gamma \stackrel{\triangle}{=} J_x J_z - J_{xz}^2$$

Rotational Dynamics

$$\dot{\boldsymbol{\omega}}_{b/i}^{b} = \mathbf{J}^{-1} \left[-\boldsymbol{\omega}_{b/i}^{b} \times \left(\mathbf{J} \boldsymbol{\omega}_{b/i}^{b} \right) + \mathbf{m}^{b} \right]$$
Define $\mathbf{m}^{b} \stackrel{\triangle}{=} \begin{pmatrix} l \\ m \\ n \end{pmatrix}$ $\mathbf{a} \times \mathbf{b} = \begin{pmatrix} 0 & -a_{3} & a_{2} \\ a_{3} & 0 & -a_{1} \\ -a_{2} & a_{1} & 0 \end{pmatrix} \begin{pmatrix} b_{1} \\ b_{2} \\ b_{3} \end{pmatrix}$

$$\begin{split} \begin{pmatrix} \dot{p} \\ \dot{q} \\ \dot{r} \end{pmatrix} &= \begin{pmatrix} \frac{J_z}{\Gamma} & 0 & \frac{J_{xz}}{\Gamma} \\ 0 & \frac{1}{J_y} & 0 \\ \frac{J_{xz}}{\Gamma} & 0 & \frac{J_x}{\Gamma} \end{pmatrix} \begin{bmatrix} \begin{pmatrix} 0 & r & -q \\ -r & 0 & p \\ q & -p & 0 \end{pmatrix} \begin{pmatrix} J_x & 0 & -J_{xz} \\ 0 & J_y & 0 \\ -J_{xz} & 0 & J_z \end{pmatrix} \begin{pmatrix} p \\ q \\ r \end{pmatrix} + \begin{pmatrix} l \\ m \\ n \end{pmatrix} \end{bmatrix} \\ &= \begin{pmatrix} \frac{J_z}{\Gamma} & 0 & \frac{J_{xz}}{\Gamma} \\ 0 & \frac{1}{J_y} & 0 \\ \frac{J_{xz}}{\Gamma} & 0 & \frac{J_x}{\Gamma} \end{pmatrix} \begin{bmatrix} \begin{pmatrix} J_{xz}pq + (J_y - J_z)qr \\ J_{xz}(r^2 - p^2) + (J_z - J_x)pr \\ (J_x - J_y)pq - J_{xz}qr \end{pmatrix} + \begin{pmatrix} l \\ m \\ n \end{pmatrix} \end{bmatrix} \\ &= \begin{pmatrix} \Gamma_1pq - \Gamma_2qr + \Gamma_3l + \Gamma_4n \\ \Gamma_5pr - \Gamma_6(p^2 - r^2) + \frac{1}{J_y}m \\ \Gamma_7pq - \Gamma_1qr + \Gamma_4l + \Gamma_8n \end{pmatrix} \end{split}$$

 Γ 's are functions of moments and products of inertia

Equations of Motion

$$\begin{pmatrix} \dot{p}_n \\ \dot{p}_e \\ \dot{p}_d \end{pmatrix} = \begin{pmatrix} c_\theta c_\psi & s_\phi s_\theta c_\psi - c_\phi s_\psi & c_\phi s_\theta c_\psi + s_\phi s_\psi \\ c_\theta s_\psi & s_\phi s_\theta s_\psi + c_\phi c_\psi & c_\phi s_\theta s_\psi - s_\phi c_\psi \\ -s_\theta & s_\phi c_\theta & c_\phi c_\theta \end{pmatrix} \begin{pmatrix} u \\ v \\ v \\ -s_\theta & s_\phi c_\theta & c_\phi c_\theta \end{pmatrix} \\ \begin{pmatrix} \dot{u} \\ \dot{v} \\ \dot{w} \end{pmatrix} = \begin{pmatrix} rv - qw \\ pw - ru \\ qu - pv \end{pmatrix} + \frac{1}{m} \begin{pmatrix} f_x \\ f_y \\ f_z \end{pmatrix},$$

$$\begin{pmatrix} \dot{\phi} \\ \dot{\theta} \\ \dot{\psi} \end{pmatrix} = \begin{pmatrix} 1 & \sin\phi \tan\theta & \cos\phi \tan\theta \\ 0 & \cos\phi & -\sin\phi \\ 0 & \frac{\sin\phi}{\cos\theta} & \frac{\cos\phi}{\cos\theta} \end{pmatrix} \begin{pmatrix} p \\ q \\ r \end{pmatrix} \\ \begin{pmatrix} \dot{p} \\ \dot{q} \\ \dot{r} \end{pmatrix} = \begin{pmatrix} \Gamma_1 pq - \Gamma_2 qr \\ \Gamma_5 pr - \Gamma_6 (p^2 - r^2) \\ \Gamma_7 pq - \Gamma_1 qr \end{pmatrix} + \begin{pmatrix} \Gamma_3 l + \Gamma_4 n \\ \frac{1}{J_y} m \\ \Gamma_4 l + \Gamma_8 n \end{pmatrix}$$

External Forces and Moments

$$\mathbf{f} = \mathbf{f}_g + \mathbf{f}_a + \mathbf{f}_p$$

 $\mathbf{m} = \mathbf{m}_a + \mathbf{m}_p$

gravitational, aerodynamic, propulsion

Gravity Force

$$\mathbf{f}_g^v = \begin{pmatrix} 0\\ 0\\ \mathsf{m}g \end{pmatrix}$$

expressed in vehicle frame

$$\mathbf{f}_{g}^{b} = \mathcal{R}_{v}^{b} \begin{pmatrix} 0\\0\\\mathsf{m}g \end{pmatrix}$$
$$= \begin{pmatrix} -\mathsf{m}g\sin\theta\\\mathsf{m}g\cos\theta\sin\phi\\\mathsf{m}g\cos\theta\cos\phi \end{pmatrix}$$

expressed in body frame